

MATH1130: Calculus II

V.2. APPLICATIONS OF LINE INTEGRALS OF VECTOR FIELDS (Not examinable!)

Definition. A vector field $\mathbf{F} : \mathbb{R}^n \rightarrow \mathbb{R}^n$ is said to be *conservative* if

$$\int_{C_1} \langle \mathbf{F}, d\mathbf{r} \rangle = \int_{C_2} \langle \mathbf{F}, d\mathbf{r} \rangle$$

for any two curves C_1 and C_2 in \mathbb{R}^n connecting the same two points \mathbf{x}_0 and \mathbf{x}_1 .

Definition. A vector field $\mathbf{F} : \Omega \rightarrow \mathbb{R}^3$ is said to be *irrotational* if $\text{curl } \mathbf{F} = \mathbf{0}$.

Proposition V.2.2. *In view of the equation $\text{curl}(\text{grad } f) = \mathbf{0}$, the following proposition can be stated. Let $\mathbf{F} : D \rightarrow \mathbb{R}^3$ be a vector field (where $D \subset \mathbb{R}^3$). Then the following three statements are equivalent:*

- \mathbf{F} is irrotational;
- there exists a scalar field $\phi : D \rightarrow \mathbb{R}$ such that $\mathbf{F} = \nabla\phi$ (ϕ is called a scalar potential of \mathbf{F});
- \mathbf{F} is conservative (that is, the work integral is independent of the path).

Calculating Scalar Potentials. Let $\mathbf{F} : D \rightarrow \mathbb{R}^3$ be a conservative vector field and let C be an arbitrary curve from a point $\mathbf{a} \in D \subset \mathbb{R}^3$ to \mathbf{x} . Then

$$\phi(\mathbf{x}) = \int_C \langle \mathbf{F}, d\mathbf{r} \rangle$$

is a scalar potential for \mathbf{F} .

For work integrals along general curves the Fundamental Theorem of Calculus takes the following form. Let $\phi : D \rightarrow \mathbb{R}$ be a continuously differentiable scalar field and let C be a curve in D from \mathbf{x}_0 to \mathbf{x}_e . Then

$$\int_C \langle \nabla\phi, d\mathbf{r} \rangle = \phi(\mathbf{x}_e) - \phi(\mathbf{x}_0).$$

Definition. A vector field $\mathbf{F} : D \rightarrow \mathbb{R}^3$ is called *solenoidal* (or *divergence-free*) if $\text{div } \mathbf{F} = 0$.

Application to Electrostatics. Given an electric field \mathbf{E} and no sources, Maxwell's equations state that

$$\text{div } \mathbf{E} = 0, \quad \text{curl } \mathbf{E} = \mathbf{0}.$$

Now

$$\text{curl } \mathbf{E} = \mathbf{0} \iff \exists \phi : D \rightarrow \mathbb{R} \text{ such that } \mathbf{E} = \nabla\phi$$

and

$$\text{div } \mathbf{E} = 0 \iff \langle \nabla, \nabla\phi \rangle = 0 \iff \nabla^2\phi = 0.$$

Therefore the potential ϕ of \mathbf{E} satisfies Laplace's equation.

Please turn over!

Application to Particle Motion. Suppose $\mathbf{F} : \mathbb{R}^3 \rightarrow \mathbb{R}^3$ is a force field and that a particle of mass m moves along a curve C through this field. Further suppose that $\mathbf{r}(t)$ is the position of the particle at time $t \in [t_0, t_e]$ ($\mathbf{r}(t)$ is a parametric representation of C where the parameter is time). Recall that

- The particle velocity is $\mathbf{v}(t) = \frac{d\mathbf{r}}{dt}$;
- The kinetic energy is $K(t) = \frac{1}{2}m \cdot \|\mathbf{v}(t)\|^2$;
- The particle acceleration is $\mathbf{a}(t) = \frac{d^2\mathbf{r}}{dt^2}$;
- Newton's Second Law states that

$$\mathbf{F}(\mathbf{r}(t)) = m \frac{d^2\mathbf{r}}{dt^2}.$$

Therefore,

$$\begin{aligned} \int_C \langle \mathbf{F}, d\mathbf{r} \rangle &= \int_{t_0}^{t_e} \langle \mathbf{F}(\mathbf{r}(t)), \frac{d\mathbf{r}}{dt} \rangle dt \\ &= \int_{t_0}^{t_e} m \left\langle \frac{d^2\mathbf{r}}{dt^2}, \frac{d\mathbf{r}}{dt} \right\rangle dt \\ &= \frac{1}{2}m \int_{t_0}^{t_e} \frac{d}{dt} \left\langle \frac{d\mathbf{r}}{dt}, \frac{d\mathbf{r}}{dt} \right\rangle dt \\ &= \frac{1}{2}m \int_{t_0}^{t_e} \frac{d}{dt} \langle \mathbf{v}(t), \mathbf{v}(t) \rangle dt, \end{aligned}$$

so using the Fundamental Theorem of Calculus yields

$$\int_C \langle \mathbf{F}, d\mathbf{r} \rangle = \frac{1}{2}m \|\mathbf{v}(t_e)\|^2 - \frac{1}{2}m \|\mathbf{v}(t_0)\|^2.$$

Hence the work done is the change in kinetic energy from time t_0 to time t_e – the conservation of energy principle.